

Num. PDE 1

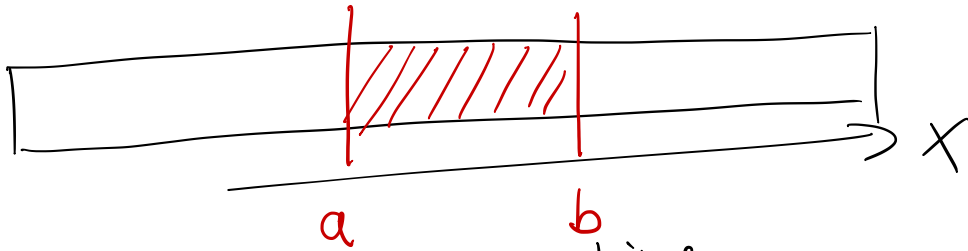
Finite volume + ERROR est.  
Convergence / consistency

Next few weeks numerical solution  
of PDE.

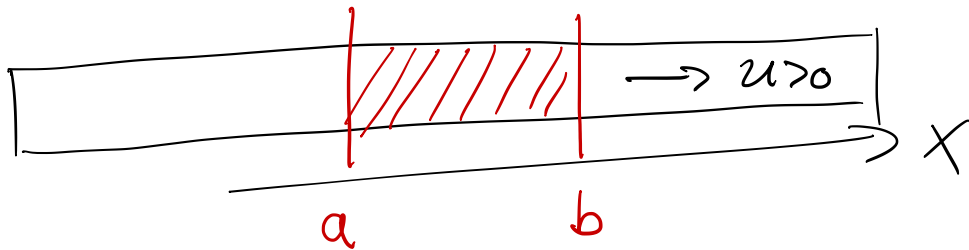
- Basic Schemes FV then DG
- Stability
- Consistency
- Convergence

Notes based on Tim Warburton's  
class notes from Math 578 at UNM.

Consider a fluid in a pipe with density  $\rho(x,t)$  flowing with a velocity  $u(x,t)$



Rate of change <sup>in time</sup> inside the red part is computed by comparing how much fluid flows out from each end



Rate of change is

$$\frac{d}{dt} \int_a^b \rho(x,t) dx = -u(b,t) \rho(b,t) + u(a,t) \rho(a,t)$$

Rate of change  
of total mass  
in section

Flux Out of  
section

Flux into section

Note that:

$$-u(b,t)g(b,t) + u(a,t)g(a,t) =$$

$$- \frac{d}{dx} \int_a^b u(x,t)g(x,t) dx \quad \text{so } \nabla$$

$$\int_a^b \frac{\partial}{\partial t} g(x,t) + \frac{\partial}{\partial x} (u(x,t)g(x,t)) dx = 0$$

*holds for any  $a, b$*   


$$\frac{\partial}{\partial t} g(x,t) + \frac{\partial}{\partial x} (g u) = 0$$

A particularly simple case is when  $u$  is constant  $u(x,t) = \bar{u}$ , then

$$\frac{\partial g}{\partial t} + u \frac{\partial g}{\partial x} = 0.$$

We can solve this by change of variables

$$\tilde{t} = t, \quad \tilde{x} = x - \bar{u}t, \quad \Rightarrow \quad \frac{\partial}{\partial t} = \frac{\partial \tilde{t}}{\partial t} \frac{\partial}{\partial \tilde{t}} + \frac{\partial \tilde{x}}{\partial t} \frac{\partial}{\partial \tilde{x}}$$

$$\begin{aligned} \text{Also: } \frac{\partial}{\partial x} &= \frac{\partial \tilde{t}}{\partial x} \frac{\partial}{\partial \tilde{t}} + \frac{\partial \tilde{x}}{\partial x} \frac{\partial}{\partial \tilde{x}} \\ &= \frac{\partial}{\partial \tilde{x}} \end{aligned}$$

$$= \frac{\partial}{\partial \tilde{t}} - \bar{u} \frac{\partial}{\partial \tilde{x}}$$

$$\underline{So:} \quad 0 = \frac{\partial}{\partial t} \rho + \bar{v} \frac{\partial}{\partial x} \rho = \left( \frac{\partial}{\partial t} - \bar{v} \frac{\partial}{\partial x} \right) \rho + \bar{v} \frac{\partial}{\partial x} \rho$$

$$= \frac{\partial}{\partial t} \rho$$

$$\frac{\partial}{\partial t} \rho = 0$$

EASY to solve

$$\rho(x, t) = \rho_0(\bar{x}) = \rho(x - \bar{v}t)$$

The solution is translated in space as time increases

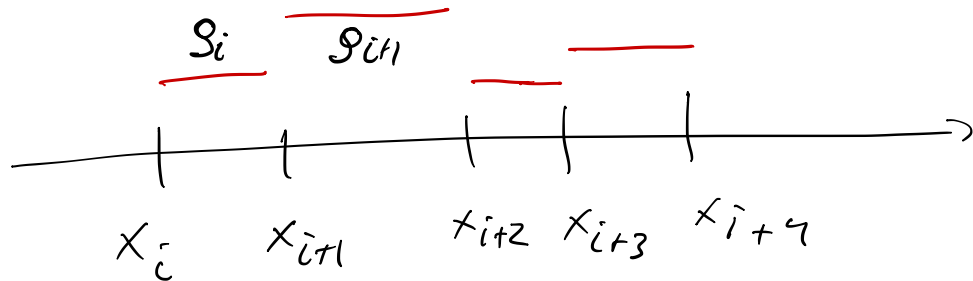


A finite volume solver from:

$$\frac{d}{dt} \int_a^b \varrho(x,t) dx = -\bar{u} \varrho(b,t) + \bar{u} \varrho(a,t)$$

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divide the real line into elements



In each element approx.  $\varrho$  by a constant  $\varrho_i$

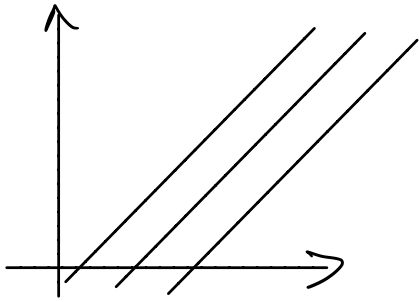


$$\frac{d}{dt} \int_{x_i}^{x_{i+1}} \rho(x,t) dx = -\bar{v} \rho(x_{i+1}, t) + \bar{v} \rho(x_i, t)$$

Note that  $\int_{x_i}^{x_{i+1}} \rho(x,t) dx = \bar{\rho}_i \cdot \Delta x$  so we can approx.

$$\frac{1}{\Delta t} (\bar{\rho}_i(t_{n+1}) - \bar{\rho}_i(t_n)) \cdot \Delta x = -\bar{v} \rho(x_{i+1}, t_n) + \bar{v} \rho(x_i, t_n)$$

How to recover values at  $x_i, x_{i+1}$   
From averages?



Use the upwind values  
to evaluate the values at  
the boundaries. Here <sup>upwind</sup>  
is to the left.

$$\frac{1}{\Delta t} (\bar{\vartheta}_i(t_{n+1}) - \bar{\vartheta}_i(t_n)) \cdot \Delta x = -\bar{u} \vartheta(x_{i+1}, t_n) + \bar{u} \vartheta(x_i, t_n)$$

$$= -\bar{u} \bar{\vartheta}_i(t_n) + \bar{u} \bar{\vartheta}_{i-1}(t_n)$$

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Simplify:  $\bar{\vartheta}_i^{n+1} = \left(1 - \bar{u} \frac{\Delta t}{\Delta x}\right) \bar{\vartheta}_i^n + \left(\bar{u} \frac{\Delta t}{\Delta x}\right) \bar{\vartheta}_{i-1}^n$

$$\bar{\vartheta}_i^{n+1} = (1 - \lambda) \bar{\vartheta}_i^n + \lambda \bar{\vartheta}_{i-1}^n \quad i = 1, 2, \dots$$

$$\vartheta_0^n = \text{Boundary cond.}$$

We now have a reasonable scheme but how good is it?

Let  $q$  be the exact solution and

$$\bar{q}_i^n = \frac{1}{x_{i+1} - x_i} \int_{x_i}^{x_{i+1}} q(x, n \cdot \Delta t) dx = \frac{1}{\Delta x} \int_{x_i}^{x_{i+1}} q(x, n \Delta t) dx$$

$q$  satisfies

$$\frac{d}{dt} \int_{x_i} q(x, t) dx = \frac{d}{dt} (\Delta x \bar{q}_i) = -\bar{u} q(x_{i+1}, t) + \bar{v} q(x_i, t)$$

We want to estimate the error, i.e. the difference between the numerical and exact solution at time  $T = n \Delta t$ ,

- We thus seek the error  $E_i^n = |\bar{q}_i^n - \bar{q}_i|$ ,  $n = \frac{T}{\Delta t}$ .
- To take limits we insist that  $\Delta t$  and  $\Delta x$  are coupled by  $\Delta t = C \cdot \Delta x$  for some fixed  $C$ .
- If we take  $\Delta t \rightarrow 0 \Rightarrow E^n = \mathcal{O}(\Delta t^5)$   
we say that the method is 5th order accurate.

$$\text{Let } \|E\|_p = \left( \Delta t \sum |E_i|^p \right)^{1/p}.$$

The scheme is said to be convergent in the norm  $\|\cdot\|$  at time  $T$  if

$$\lim_{\Delta t \rightarrow 0} \|E^n\| = 0$$

$$\Delta t \rightarrow 0$$

$$T = n \Delta t$$

Plan of attack, Reduce errors at subsequent timesteps and find a bound at  $T$ .

We get to use the method

We start by estimating the local truncation error.

The local truncation error is the error committed in one timestep assuming we start with exact initial data

$$\therefore \bar{g}_i^n = \bar{q}_i^n, \text{ and } \bar{g}_i^{n+1} = \left(1 - \frac{\Delta t}{\Delta x} \bar{u}\right) \bar{q}_i^n + \frac{\Delta t}{\Delta x} \bar{u} \bar{q}_{i-1}^n$$

$$\text{LTE} = R_i^n \equiv \frac{1}{\Delta t} (\bar{g}_i^{n+1} - \bar{q}_i^{n+1})$$

TAYLOR

$$\underline{q}_{i-1}^n = \bar{q}_i^n - \Delta x \left( \frac{\partial \bar{q}}{\partial x} \right) + \frac{\Delta x^2}{2} \left( \frac{\partial^2 \bar{q}}{\partial x^2} \right) + \mathcal{O}(\Delta x^3)$$

$$\underline{q}_{i-1}^{n+1} = \bar{q}_i^n + \Delta t \left( \frac{\partial \bar{q}}{\partial t} \right) + \frac{\Delta t^2}{2} \left( \frac{\partial^2 \bar{q}}{\partial t^2} \right) + \mathcal{O}(\Delta t^3)$$

$$R_i^n = \frac{1}{\Delta t} \left( \left( 1 - \frac{\Delta t}{\Delta x} \bar{u} \right) \bar{q}_i^n + \frac{\Delta t}{\Delta x} \bar{u} \underline{q}_{i-1}^n - \underline{q}_{i-1}^{n+1} \right) \quad \text{Simplify}$$

DO IT!

$$= - \left( \frac{\partial \bar{q}}{\partial t} + \bar{u} \frac{\partial \bar{q}}{\partial x} \right) - \left\{ \frac{\Delta t}{2} \left( \frac{\partial^2 \bar{q}}{\partial t^2} \right) - \frac{\bar{u} \Delta x}{2} \left( \frac{\partial^2 \bar{q}}{\partial x^2} \right) + \mathcal{O}(\Delta x^2) \right. \\ \left. + \mathcal{O}(\Delta t^2) \right\}$$

Simplify more:  $\overline{\frac{\partial q}{\partial t}} + \pi \overline{\frac{\partial q}{\partial x}} = 0$ ,  $\overline{\frac{\partial^2 q}{\partial t^2}} = \pi^2 \overline{\frac{\partial^2 u}{\partial x^2}}$

$$R_i^n = \frac{\bar{n} \Delta x}{2} \left(1 - \frac{\bar{n} \Delta t}{\Delta x}\right) \overline{\frac{\partial^2 q}{\partial x^2}} + \underline{O(\Delta x^2)},$$

Note that  $R_i^n \rightarrow 0$  as  $\Delta t \rightarrow 0$   
assuming  $\Delta t = C \cdot \Delta x$

When the local truncation error  $\rightarrow 0$  as  $\Delta t \rightarrow 0$   
we say that the method is consistent



The local truncation error predicts that we will make an  $\Delta t$  error every timestep. Let's be a bit more precise

DEF  $\bar{q}_i^n = \bar{q}_i^n + E_i^n$  and write the scheme as

$$\bar{q}^{n+1} = N \bar{q}^n.$$

We have 
$$\begin{aligned} \bar{E}^{n+1} &= N(\underbrace{\bar{q}^n + E^n}_{\bar{q}^n}) - \bar{q}^{n+1} = N(\bar{q}^n + \bar{E}^n) - N(\bar{q}^n) + N(\bar{q}^n) - \bar{q}^{n+1} \\ &= \underbrace{N(\bar{q}^n + \bar{E}^n) - N(\bar{q}^n)}_{\text{propagation of previous ERROR}} + \Delta t R^n \end{aligned}$$

TRUNCATION  
ERROR

Now suppose  $N$  is for some norm  $N$  is a contraction.

$$\|N(p) - N(q)\| \leq \|p - q\|$$

$$E^{n+1} = N(\bar{q}^n + E^n) - N(\bar{q}^n) + \Delta t R^n, \text{ taking norms + } \Delta \text{ ineq.}$$

$$\|E^{n+1}\| \leq \|N(\bar{q}^n + E^n) - N(\bar{q}^n)\| + \Delta t \|R^n\|$$

$$\leq \|\bar{q}^n + E^n - \bar{q}^n\| + \Delta t \|R^n\| \leq \|E^n\| + \Delta t \|R^n\|$$

$$\leq \|E^{n-1}\| + \Delta t \|R^{n-1}\| + \Delta t \|R^n\| \quad \left\{ \text{many times} \right\}$$

$$\leq \|E^0\| + \Delta t \sum_{m=1}^n \|R^m\|$$

In other words the error at time  $T$  is

$$\|E^{nH}\| \leq \|E^0\| + \Delta t \sum_{m=1}^n \|R^m\| \leq \|E^0\| + T \max_{m=1, \dots, n} \|R^m\|$$

So IF the method is consistent (and  $q$  is smooth)

$$\|E^{nH}\| \leq \|E^0\| + T \mathcal{O}(\Delta x)$$



Initial data  $\rightarrow 0$  as  $\Delta x \rightarrow 0$

Home work : Prove that

$\bar{g}_i^n = (1-\lambda) \bar{g}_i^n + \lambda \bar{g}_{i-1}^n$  is a contraction  
in  $\|\cdot\|_1$  when  $\bar{g}_0^n = 0$ .

# DG: Consistency - Stability - Convergence

We solved  $\frac{d}{dt} (\Delta x \bar{q}_i(t)) = -\bar{u} q(x_{i+1}, t) + \bar{u} q(x_i, t)$

$$\approx -\bar{u} \bar{q}_i + \bar{u} \bar{q}_{i-1}$$

↑  
upwind approx.

Our first DG scheme will generalize this so that we can exploit approximation by high degree polynomials

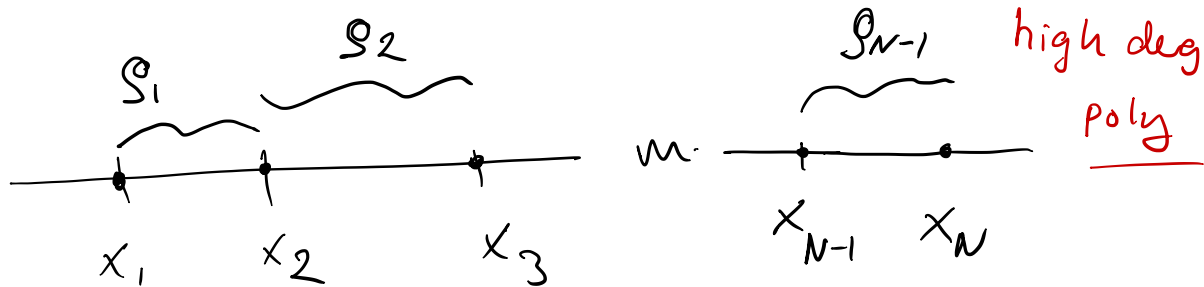
Again divide the pipe into segments separated by  $x_1, x_2, x_3, \dots, x_N$ .

On each segment  $x_i \leq x \leq x_{i+1}$  we solve a weak form of the advection eq.

Find  $g_i^s \in P^P(x_i, x_{i+1})$  such that  $\forall v \in P^P(x_i, x_{i+1})$

$$\int_{x_i}^{x_{i+1}} v \left( \frac{\partial g_i^s}{\partial t} + \bar{u} \frac{\partial g_i^s}{\partial x} \right) dx = v(x_i) \bar{u} \left( g_{i-1}^s(x_i) - g_i^s(x_i) \right)$$

*Space of polynomials of deg P*



$$\int_{x_i}^{x_{i+1}} v \left( \frac{\partial g_{i+1}}{\partial t} + \bar{u} \frac{\partial g_i}{\partial x} \right) dx = v(x_i) \bar{u} (g_{i-1}(x_i) - g_i(x_i))$$

$x_i$

Selecting  $g_i \in P^0(x_i, x_{i+1})$  i.e. constants

Reduces this to finite vol.

What about stability?

The best part of dG is that this is trivial

- Find  $g_i \in P^p(x_i, x_{i+1})$  such that  $\forall v \in P^p(x_i, x_{i+1})$

FOR ALL

$$\int_{x_i}^{x_{i+1}} v \left( \frac{\partial g_i}{\partial t} + \bar{n} \frac{\partial g_i}{\partial x} \right) dx = v(x_i) (g_{i-1}(x_i) - g_i(x_i))$$

Select  $v = g_i$ !

$$\int_{x_i}^{x_{i+1}} g_i \left( \frac{\partial g_i}{\partial t} + \bar{n} \frac{\partial g_i}{\partial x} \right) dx = g_i(x_i) (g_{i-1}(x_i) - g_i(x_i))$$

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Simplify to get an energy estimate

$$\frac{d}{dt} \int_{x_i}^{x_{i+1}} \frac{\rho_i^2}{2} dx = -\bar{u} \left( \int_{x_i}^{x_{i+1}} \rho_i \frac{\partial \rho_i}{\partial x} dx \right) + \bar{u} \rho_i(x_i) (\rho_{i-1}(x_i) - \rho_i(x_i))$$

$$\underline{\text{IBP}} = -\bar{u} \left[ \frac{\rho_i^2}{2} \right]_{x_i}^{x_{i+1}} + \bar{u} \rho_i(x_i) (\rho_{i-1}(x_i) - \rho_i(x_i))$$

$$= \bar{u} \left( -\frac{\rho_i^2(x_i)}{2} - \frac{\rho_i^2(x_{i+1})}{2} \right) + \bar{u} \rho_i(x_i) \rho_{i-1}(x_i)$$

$$\frac{d}{dt} \int_{x_i}^{x_{i+1}} \frac{\rho_i^2}{2} dx = \bar{u} \left( -\frac{\rho_i^2(x_i)}{2} - \frac{\rho_i^2(x_{i+1})}{2} \right) + \bar{u} \rho_i(x_i) \rho_{i-1}(x_i)$$

Sum over ELEMS

$$\sum_{i=1}^{N-1} : \frac{d}{dt} \int_{x_i}^{x_{i+1}} \frac{\rho_i^2}{2} dx = \bar{u} \left( -\frac{\rho_i^2(x_i)}{2} - \frac{\rho_i^2(x_{i+1})}{2} \right) + \bar{u} \rho_i(x_i) \rho_{i-1}(x_i)$$

$$\frac{d}{dt} \sum_{i=1}^{N-1} \left( \int_{x_i}^{x_{i+1}} \frac{\rho_i^2}{2} dx \right) = \sum_{i=1}^{N-1} \left[ \bar{u} \left( -\frac{\rho_i^2(x_i)}{2} - \frac{\rho_i^2(x_{i+1})}{2} \right) + \bar{u} \rho_i(x_i) \rho_{i-1}(x_i) \right]$$

$$\leq \sum_{i=1}^{N-1} \left[ \bar{u} \left( \frac{\rho_i^2(x_i)}{2} + \frac{\rho_{i-1}^2(x_i)}{2} \right) \right]$$

$$\leq \frac{\bar{u}}{2} \sum_{i=1}^{N-1} \left[ -\rho_i^2(x_{i+1}) + \rho_{i-1}^2(x_i) \right]$$

$$\leq \frac{\bar{u}}{2} \left( \rho_0^2(x_1) - \rho_{N-1}^2(x_N) \right)$$

Note that we left time continuous. If time derivatives are discretized exactly we have that the energy of the numerical solution is controlled by the boundary condition  $\rho_0(x_i)$

denoting  $\|\rho\|_2^2 = \sum_{i=1}^{N-1} \int_{x_i}^{x_{i+1}} \frac{\rho_i^2}{2} dx$  we find

$$\frac{d}{dt} \frac{\|\rho\|_2^2}{2} \leq \frac{\bar{u}}{2} \left( \rho_0^2(x_1) - \rho_{N-1}^2(x_N) \right)$$

$$\text{or } \frac{d}{dt} \|\rho\|_2 \leq \frac{\bar{u}}{2\|\rho\|_2} \left( \rho_0^2(x_1) - \rho_{N-1}^2(x_N) \right)$$

Now we can try to bound the error in the same way.

Eg.: Let  $\rho$  and  $\sigma$  be two numerical solutions that are close at the initial time, then trivially (same BC)

$$\frac{d}{dt} \|\rho - \sigma\|_2 \leq \frac{-\bar{u} \left( \rho_{N-1}(x_N) - \sigma_{N-1}(x_N) \right)^2}{2 \|\rho - \sigma\|_2}$$

$\leq 0 \rightarrow$  the solutions get closer with time.

1. Let  $q$  be the exact solution

2. Assume that we have a projection operator

$$\Pi^P : H^1(x_i, x_N) \rightarrow \bigcup_{i=1}^{N-1} P^P(x_i, x_{i+1})$$

We work with 3 - equations

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$$\textcircled{1} \int v \Pi^P \left( \frac{\partial q_i}{\partial t} + \bar{u} \frac{\partial q_i}{\partial x} \right) dx = 0 \quad \begin{array}{l} \text{TRUNCATION} \\ \text{LOC.} \end{array}$$

$$\textcircled{2} \int v \left( \frac{\partial \Pi^P q_i}{\partial t} + \bar{u} \frac{\partial \Pi^P q_i}{\partial x} \right) dx = v(x_i) \left( \Pi^P q_{i-1}(x_i) - \Pi^P q_i(x_i) \right) + \int v R_i$$

$$\textcircled{3} \int v \left( \frac{\partial g_i}{\partial t} + \bar{u} \frac{\partial g_i}{\partial x} \right) dx = v(x_i) \left( g_{i-1}(x_i) - g_i(x_i) \right)$$

- ① Project PDE. } Assume that projection and  $\frac{\partial}{\partial t}$   
 ② Truncation error. } commute and that  $q$  is smooth.  
 ③ The method. } 1 + 2 gives

$$\int_{x_i}^{x_{i+1}} v R_i dx = \int v \left( \bar{u} \frac{\partial \Pi^P q_i}{\partial x} - \bar{u} \Pi^P \frac{\partial q_i}{\partial x} \right) dx - v(x_i) \left( \Pi^P q_{i-1}(x_i) - q_{i-1}(x_i) + q_i(x_i) - \Pi^P q_i(x_i) \right)$$

Now USE  $v = R_i \in P^p + C-S + \Delta$  ineq.

$$\|R_i\|_2 \leq \underbrace{\sum_{i=1}^{N-1} \left\| \bar{u} \frac{\partial \Pi^P q_i}{\partial x} - \bar{u} \Pi^P \frac{\partial q_i}{\partial x} \right\|_2}_{C h^s \left\| \frac{d^p q}{dx^p} \right\|_2} + \sum_{i=1}^{N-1} \left( \left| \Pi^P q_{i-1}(x_i) - q_{i-1}(x_i) \right| + \left| \Pi^P q_i(x_i) - q_i(x_i) \right| \right)$$

We have an estimate for truncation error

Now we estimate  $\|\Pi^P q - g\|_2$  by ③-②.

$$\textcircled{2} \int v \left( \frac{\partial \Pi^P q_i}{\partial t} + \bar{u} \frac{\partial \Pi^P q_i}{\partial x} \right) dx = v(x_i) \left( \Pi^P q_{i-1}(x_i) - \Pi^P q_i(x_i) \right) + \int_{x_i}^{x_{i+1}} v R_i$$

$$\textcircled{3} \int v \left( \frac{\partial g_i}{\partial t} + \bar{u} \frac{\partial g_i}{\partial x} \right) dx = v(x_i) \left( g_{i-1}(x_i) - g_i(x_i) \right)$$

Subtract ③ from ② and set  $v = (\Pi^P q_i - g_i)$

$$\int_{x_i}^{x_{i+1}} (\Pi^P q_i - g_i) \left( \frac{\partial \Pi^P q_i - g_i}{\partial t} + \bar{u} \frac{\partial \Pi^P q_i - g_i}{\partial x} \right) dx =$$

$$(\Pi^P q_i(x_i) - g_i(x_i)) \left( \Pi^P q_{i-1}(x_i) - g_{i-1}(x_i) - \Pi^P q_i(x_i) + g_i(x_i) \right)$$

$$+ \int_{x_i}^{x_{i+1}} (\Pi^P q_i - g_i) R_i$$

USE the stability of Res

$$\frac{d}{dt} \|g - \sigma\|_2 \leq 0$$

$$\frac{d}{dt} \frac{\|\Pi^P q - \mathcal{Q}\|_2^2}{2} \leq \sum_{\bar{i}=1}^{N-1} \int_{x_i}^{x_{i+1}} (\Pi^P q_i - \mathcal{Q}_i) R_i dx$$

$$\leq \|\Pi^P q - \mathcal{Q}\|_2 \|R\|_2$$

Use Estimate for  $\|R\|$

$$\frac{d}{dt} \|\Pi^P q - \mathcal{Q}\|_2 \leq \|R\| \leq C \cdot h^S \left\| \frac{d^S q}{dx^S} \right\|_2$$

The numerical sol. and the projection  
grows slowly apart in time.



We integrate in time:

$$\|\Pi^P q - g\|_2 \leq \|\Pi^P q(t=0) - g(t=0)\|_2 + Ch^S T \max_{0 \leq \tau \leq T} \left\| \frac{d^P q}{dx^P}(\tau) \right\|_2$$

Finally we estimate error between exact and num.

$$\text{Sol. } \|q - g\|_2 = \|q - \Pi q + \Pi q - g\|$$

$$\leq \|q - \Pi q\|_2 + \|\Pi q(t=0) - g(t=0)\|_2 + Ch^S T \max_{0 \leq \tau \leq T} \left\| \frac{d^P q(\tau)}{dx^P} \right\|_2$$

First term is beyond our control

Second term suggests we approximate initial data accurately.

Third term leads to linear growth with  $T$ .

- We swept  $\sin$   $Ch^S \left\| \frac{d^p q}{dx^p} \right\|_2$  under the rug
- In multi D the arguments are the same
- In special cases one can cook up a better  $\Pi$  to get a sharper estimate (Does not affect the method though)

Next lecture we start thinking about multiple dimensions.